

Prospects for understanding dark matter-baryon interactions with direct detection

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Main question:

Could we discern details of underlying DM physics, such as the type of scattering interaction, with Generation 2 of direct detection?

EFT of DM-nucleon interactions

EFT = experiment-driven way of parameterizing underlying theories.

$$V_{\text{eff}} = h_1 \delta^3(\vec{r}) - h_2 \vec{s}_\chi \cdot \vec{\nabla} \delta^3(\vec{r}) + \ell_1 \frac{1}{4\pi r} + \ell_2 \frac{\vec{s}_\chi \cdot \vec{r}}{4\pi r^3}$$

The diagram shows two purple arrows pointing from the terms $h_1 \delta^3(\vec{r})$ and $-h_2 \vec{s}_\chi \cdot \vec{\nabla} \delta^3(\vec{r})$ to the label "heavy mediator". Two green arrows point from the terms $\ell_1 \frac{1}{4\pi r}$ and $\ell_2 \frac{\vec{s}_\chi \cdot \vec{r}}{4\pi r^3}$ to the label "light mediator".

Fan et al, 2010

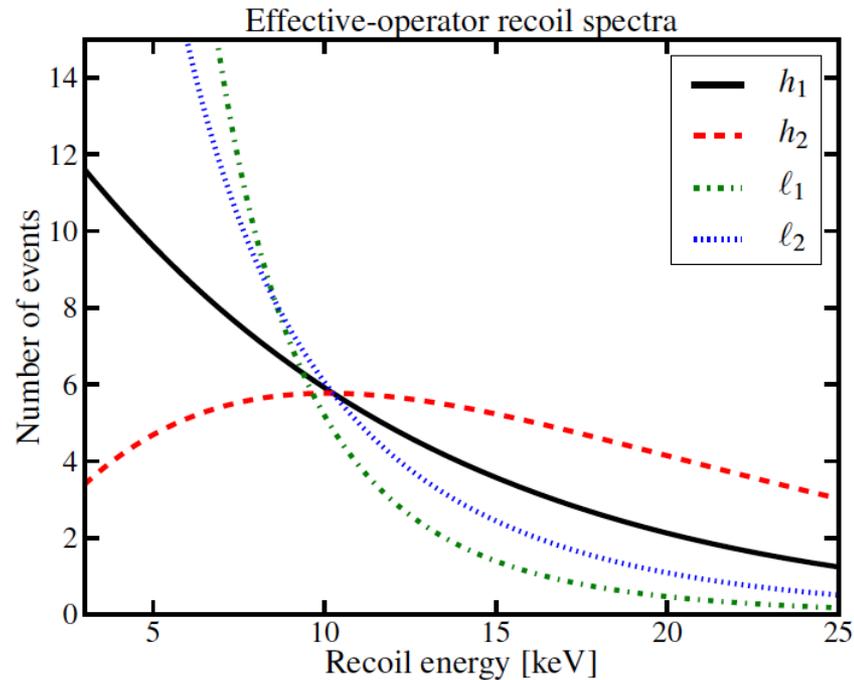
(also, McDermott et al, 2012; Feldstein et al, 2012; Chang et al, 2010; Fitzpatrick et al 2013, etc.)

EFT of DM-nucleon interactions

$$\frac{d\sigma}{dE_R} = \frac{A^2 F^2(E_R) m_N}{2\pi v^2} \left(\left| h_1 + \frac{\ell_1}{2m_N E_R} \right|^2 + \frac{1}{4} \left| h_2 \sqrt{2m_N E_R} + \frac{\ell_2}{\sqrt{2m_N E_R}} \right|^2 \right)$$

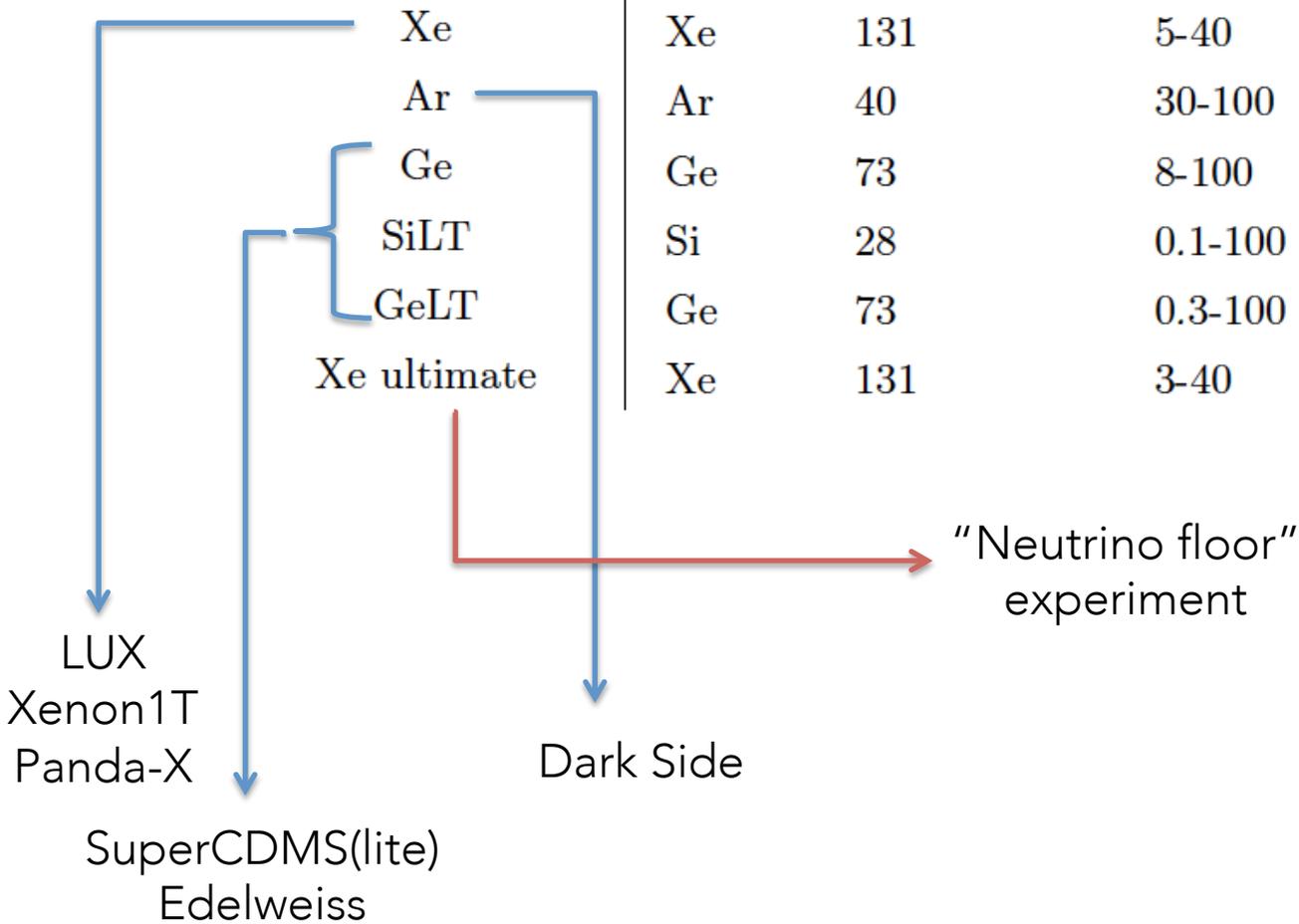
Fan et al, 2010

$$(\sigma_{SI} = \frac{\mu_p^2}{\pi} h_1^2)$$

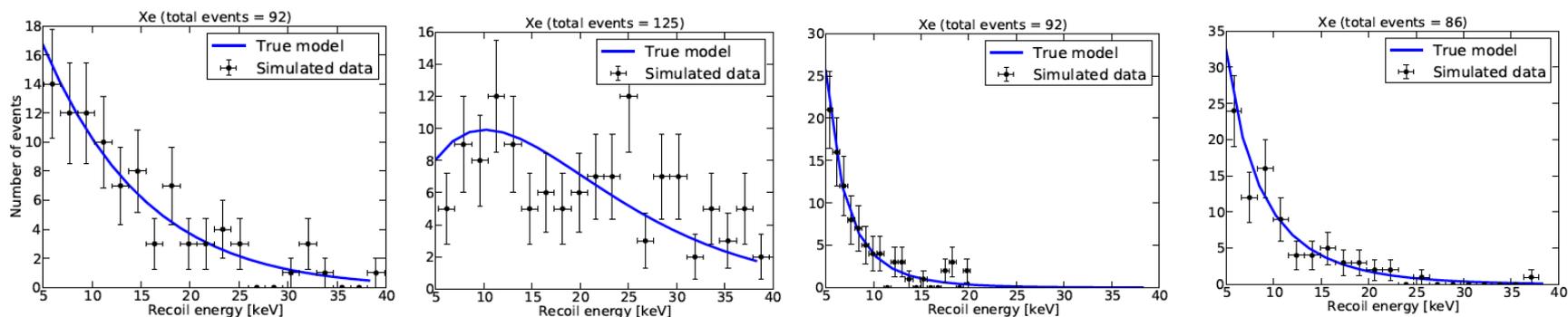


Experimental context: Generation 2

Experiment label	Target	Nuclear mass [AMU]	Energy window [keV]	Exposure [kg-yr]
Xe	Xe	131	5-40	2000
Ar	Ar	40	30-100	1000
Ge	Ge	73	8-100	100
SiLT	Si	28	0.1-100	0.2
GeLT	Ge	73	0.3-100	4
Xe ultimate	Xe	131	3-40	10 000



Method



Step 1: Simulate data ($\{E_R\}$ spectrum) + Poisson noise for a single operator with max coupling.

Step 2: Analyze each simulation by fitting each of the 4 operator models.

Step 3: Evaluate probability of the right model compared to other models.

Step 4: Repeat 1-3 for each of the 5 mock experiments, many times.

Step 5: Jointly analyze different experiments.

Framework: Bayesian model selection

coupling + mass

Posterior probability:

$$\mathcal{P}(\Theta|\{E_R\}, M) = \frac{\mathcal{L}(\{E_R\}|\Theta, M)p(\Theta|M)}{\mathcal{E}(M)}$$

Likelihood:

$$\mathcal{L}(\{E_R\}|\Theta, M) = P(N|\Theta, M) \prod_{i=1}^N P_1(E_R^i|\Theta, M)$$

Evidence:

$$\mathcal{E}(M) = \int d\Theta \mathcal{L}(\{E_R\}|\Theta, M)p(\Theta|M)$$

$$p(M_j) = \frac{\mathcal{E}(M_j)}{\sum_i \mathcal{E}(M_i)}$$

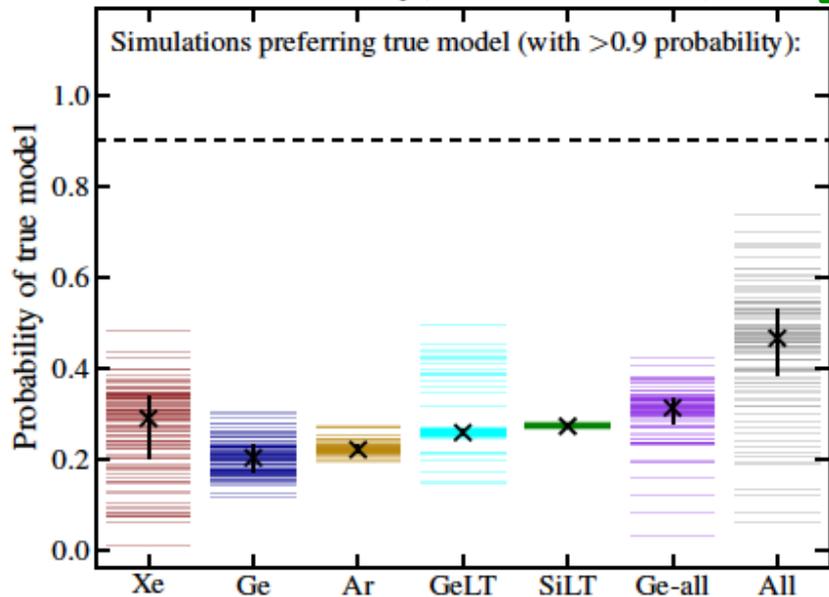


Model probability
(given data)

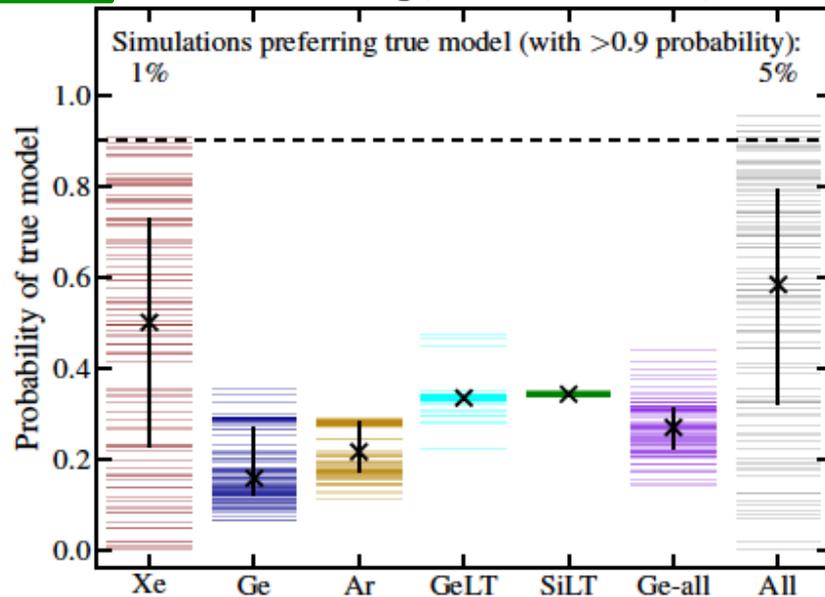
Success probability = percent of simulations in which the right operator was selected with >90% probability.

20 GeV

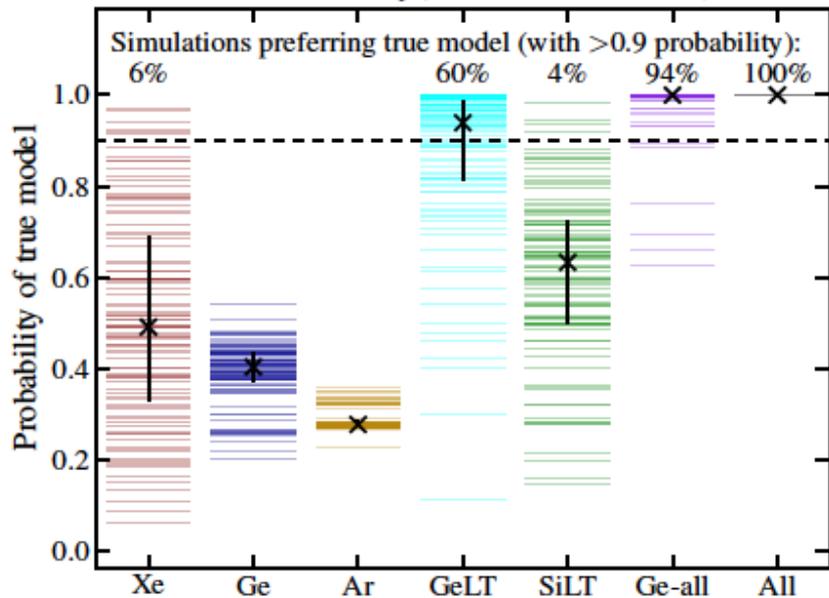
True model: h_1 (WIMP mass=20GeV)



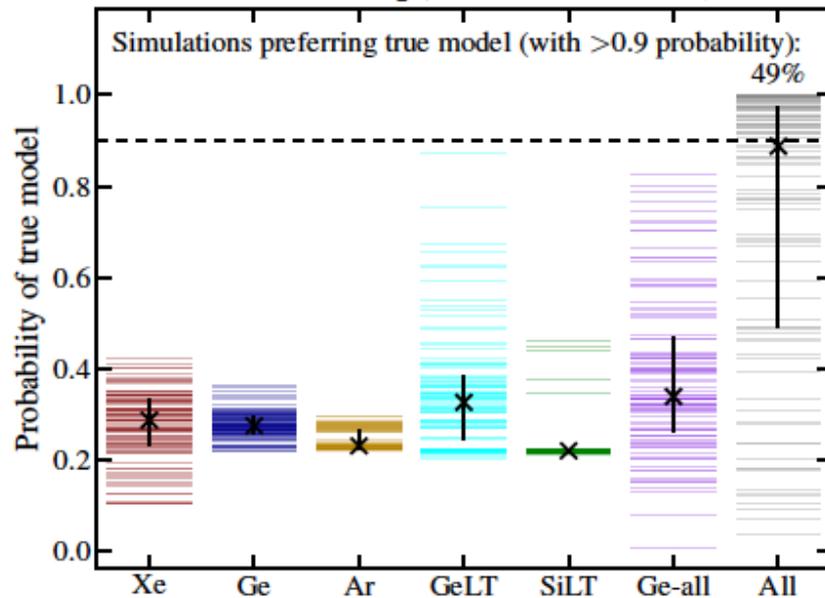
True model: h_2 (WIMP mass=20GeV)



True model: ℓ_1 (WIMP mass=20GeV)

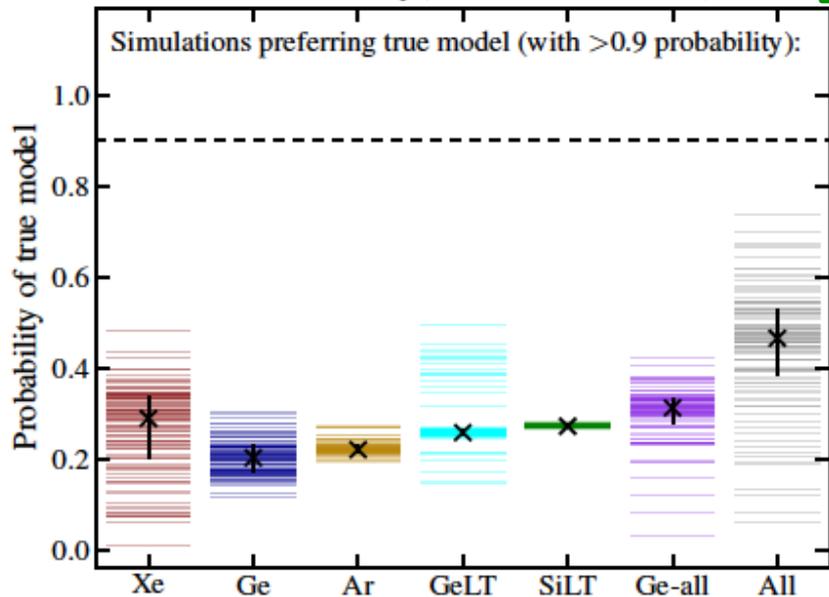


True model: ℓ_2 (WIMP mass=20GeV)

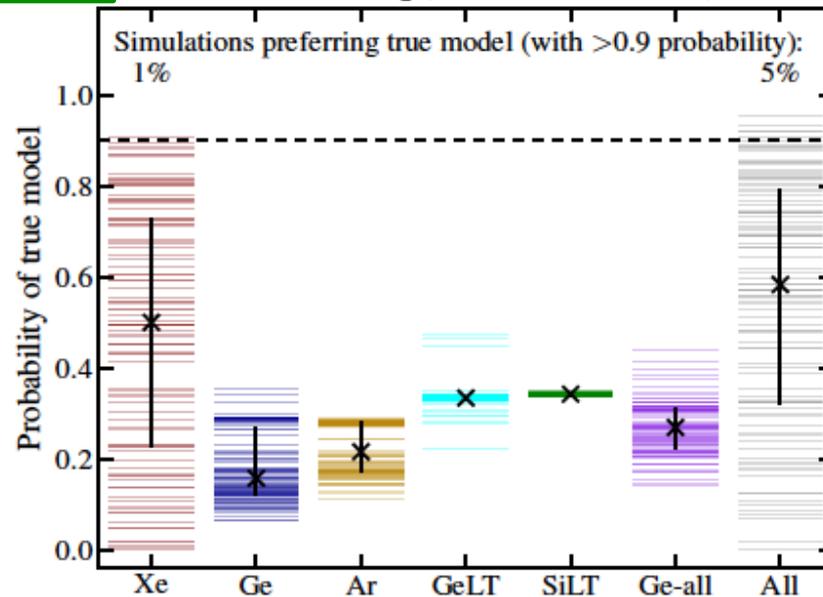


20 GeV

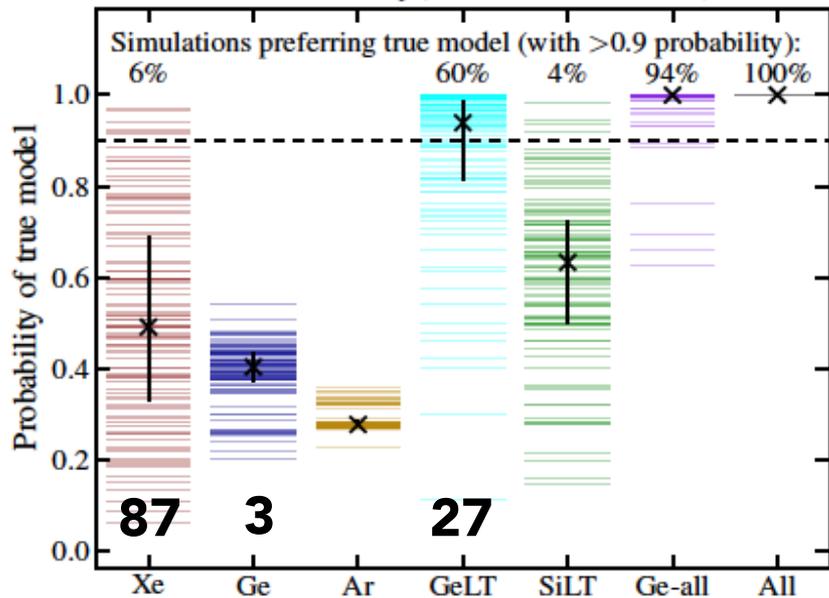
True model: h_1 (WIMP mass=20GeV)



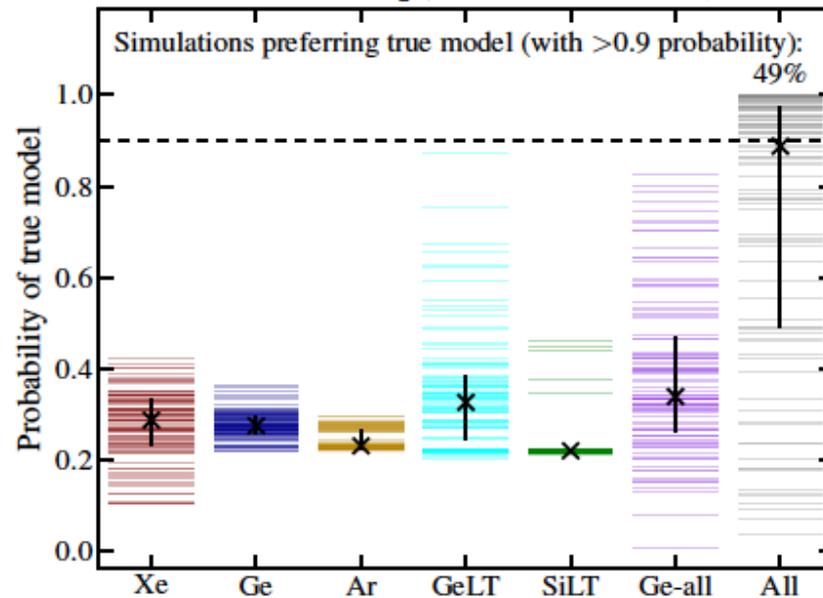
True model: h_2 (WIMP mass=20GeV)



True model: ℓ_1 (WIMP mass=20GeV)

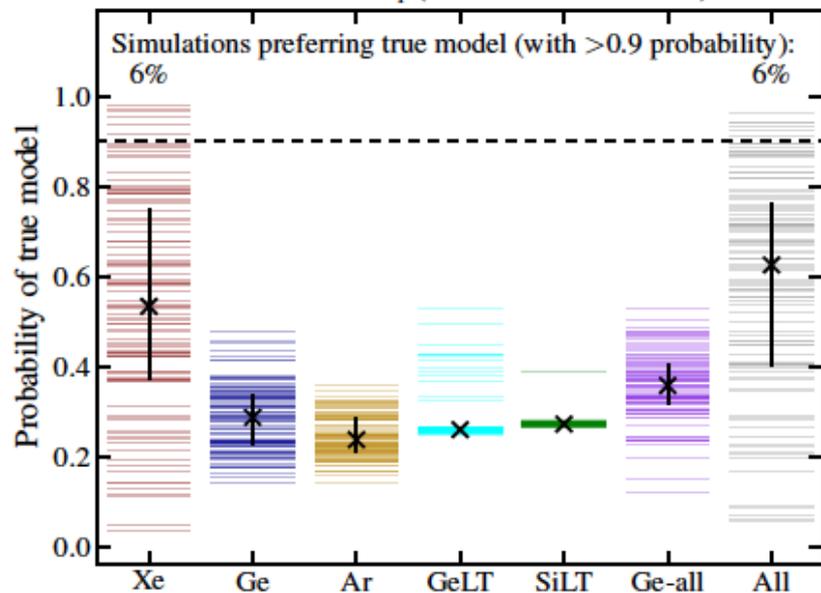


True model: ℓ_2 (WIMP mass=20GeV)

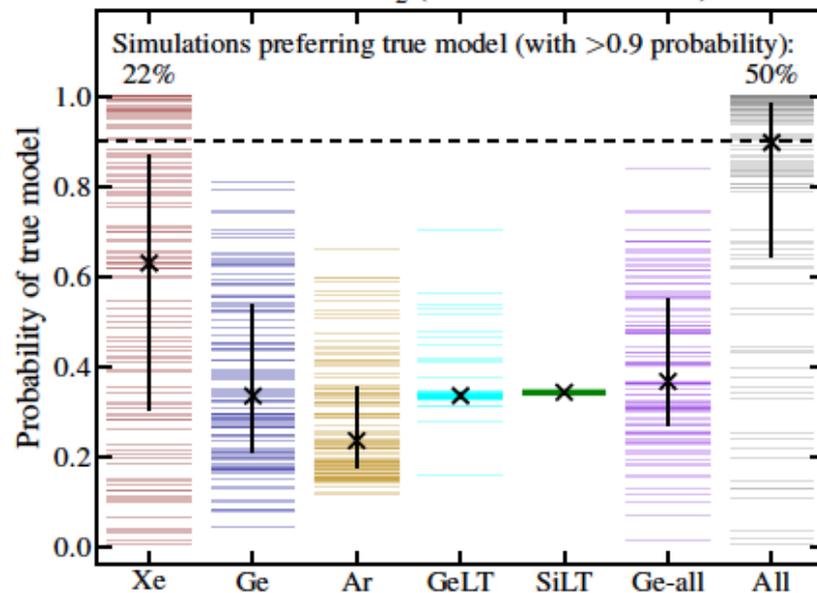


50 GeV

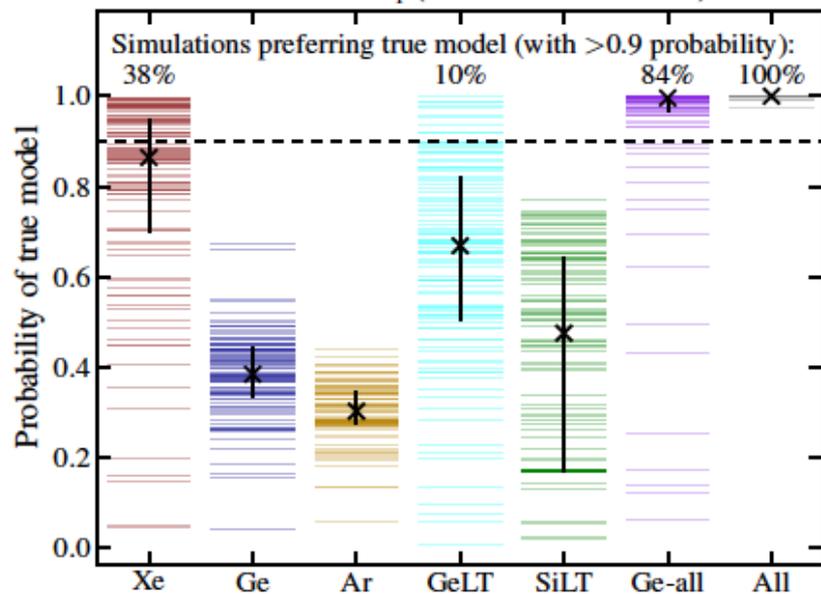
True model: h_1 (WIMP mass=50GeV)



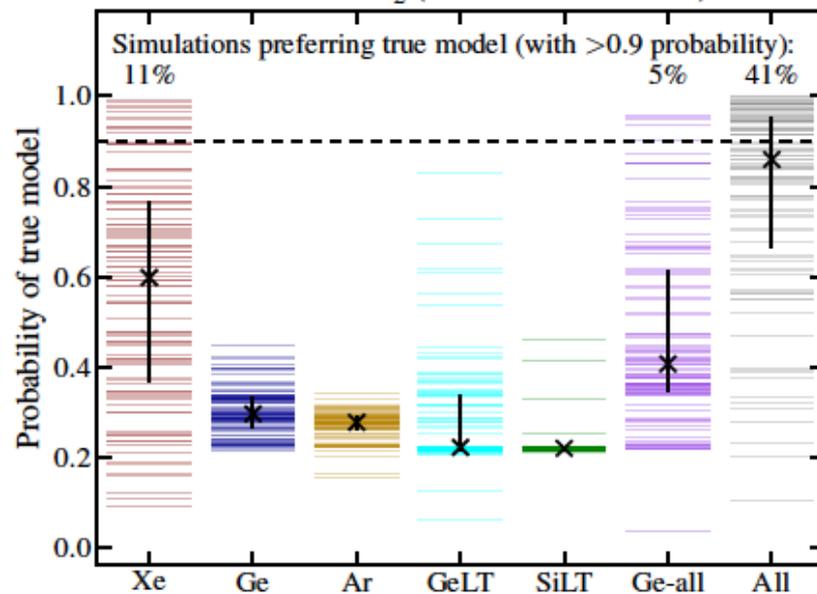
True model: h_2 (WIMP mass=50GeV)



True model: ℓ_1 (WIMP mass=50GeV)

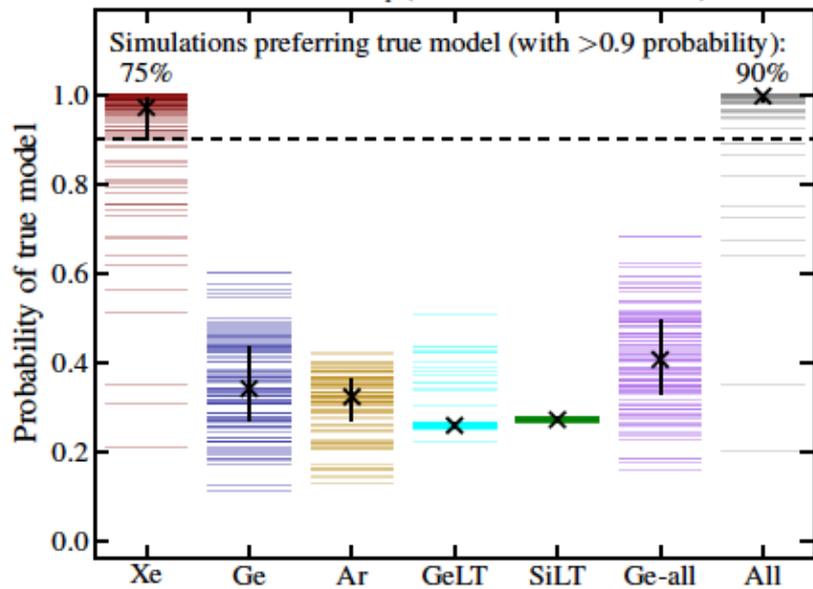


True model: ℓ_2 (WIMP mass=50GeV)

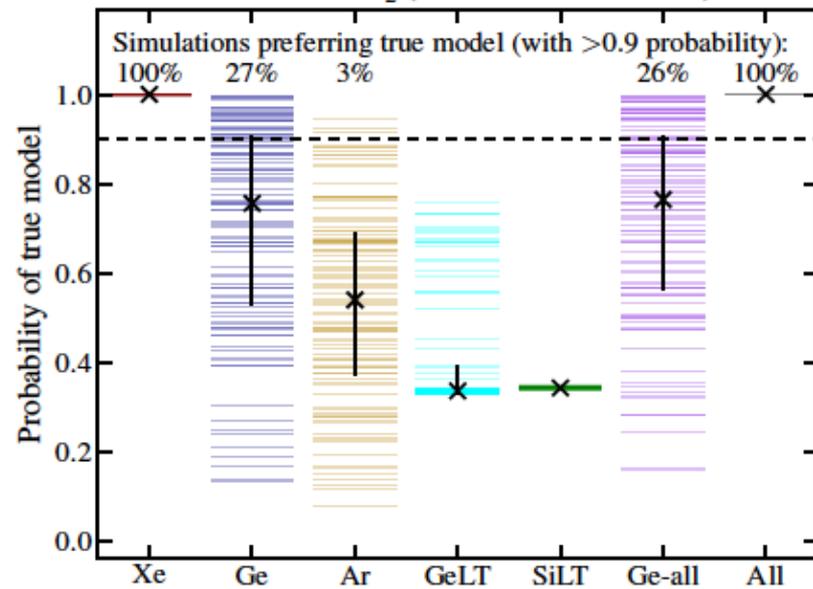


200 GeV

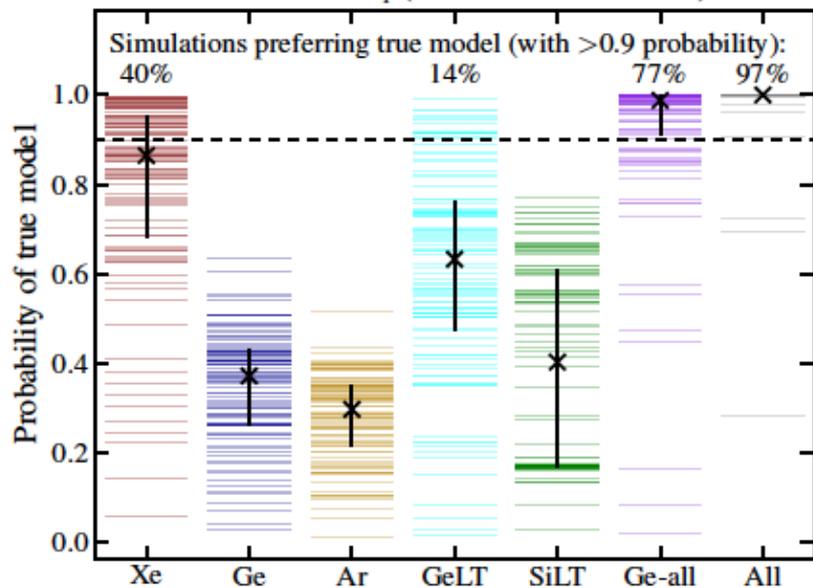
True model: h_1 (WIMP mass=200GeV)



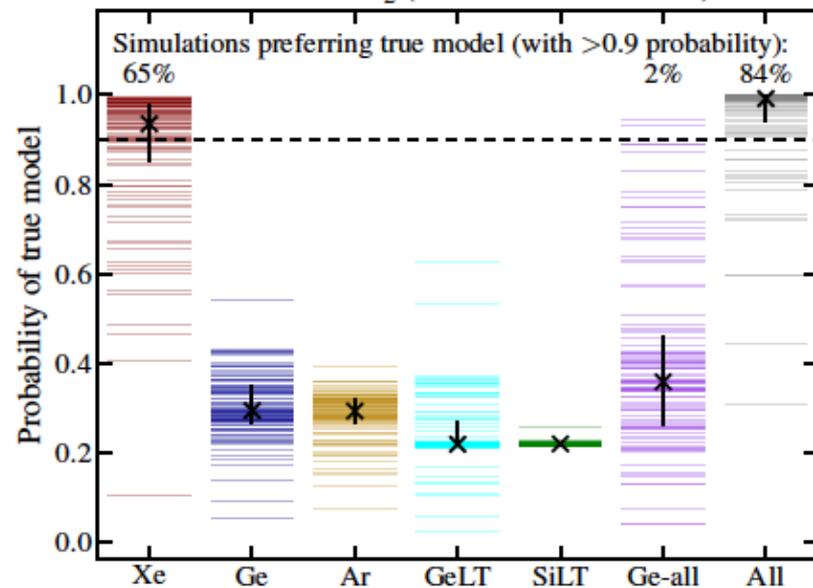
True model: h_2 (WIMP mass=200GeV)



True model: ℓ_1 (WIMP mass=200GeV)

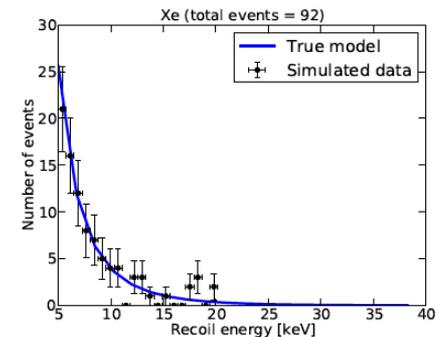
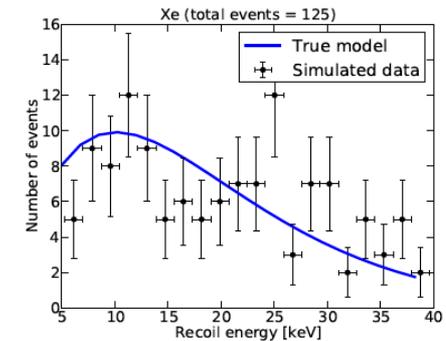
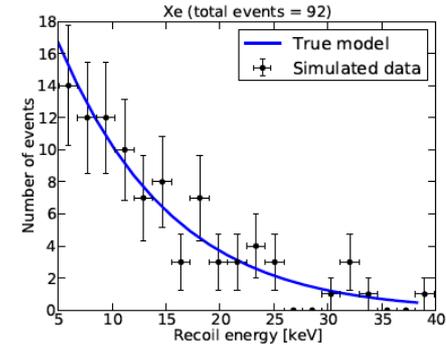


True model: ℓ_2 (WIMP mass=200GeV)

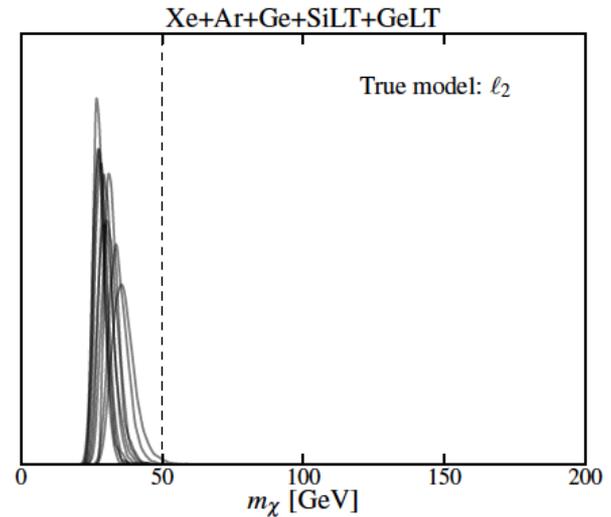
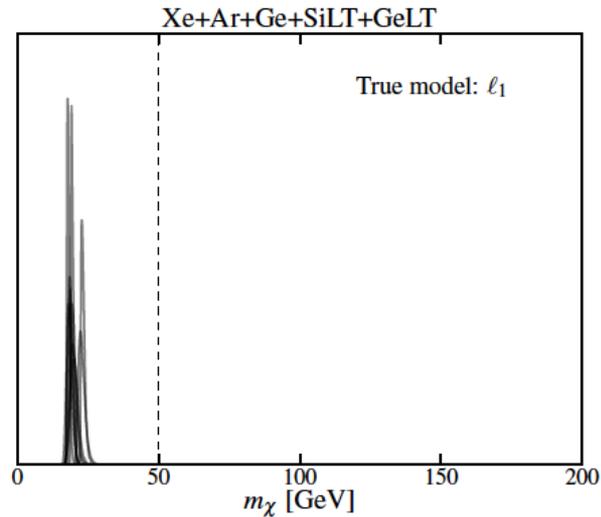
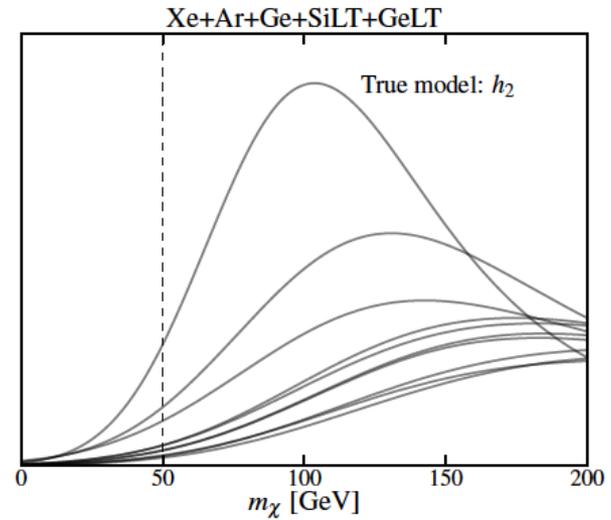
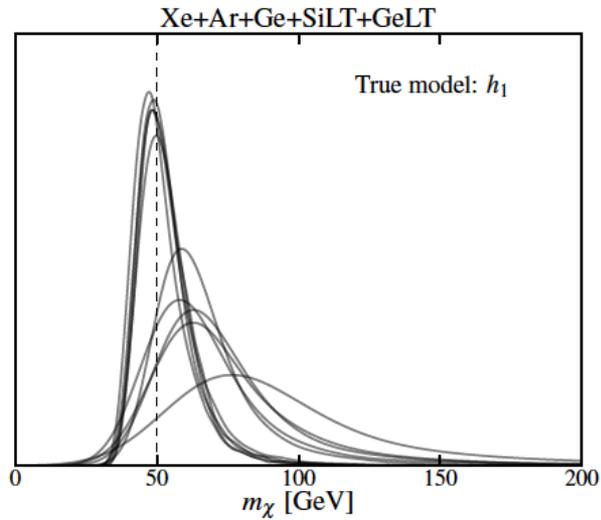


For successful model selection

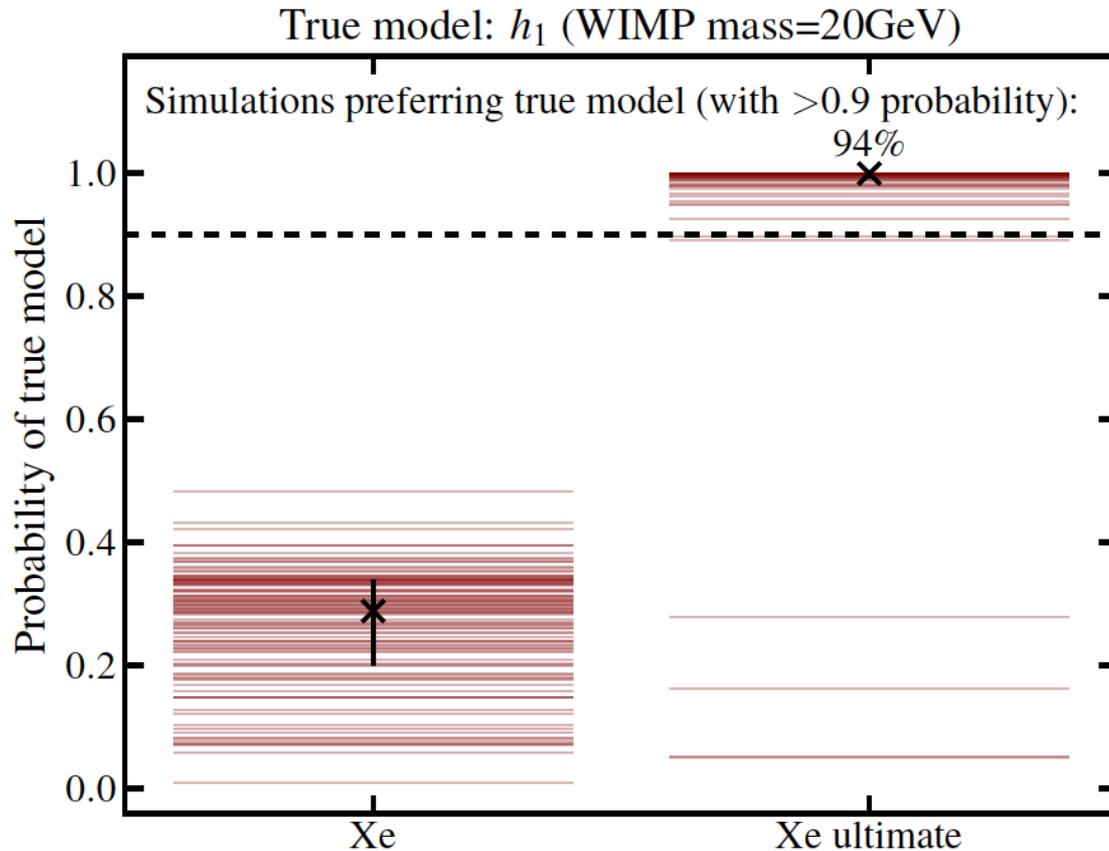
- Wide energy coverage
- Low threshold
- Large exposure
- Multiple targets



WIMP mass reconstruction: beware of bias



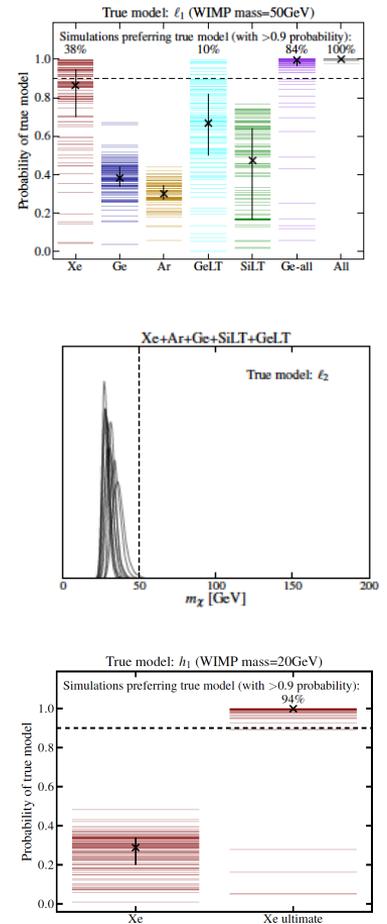
How far can model selection go for a signal at the current limit?



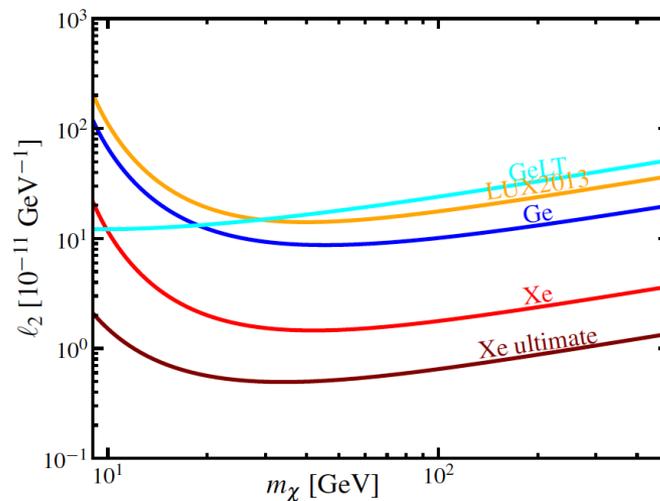
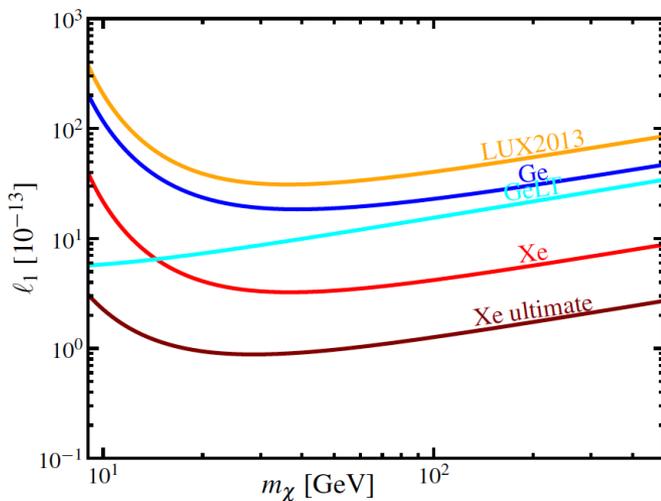
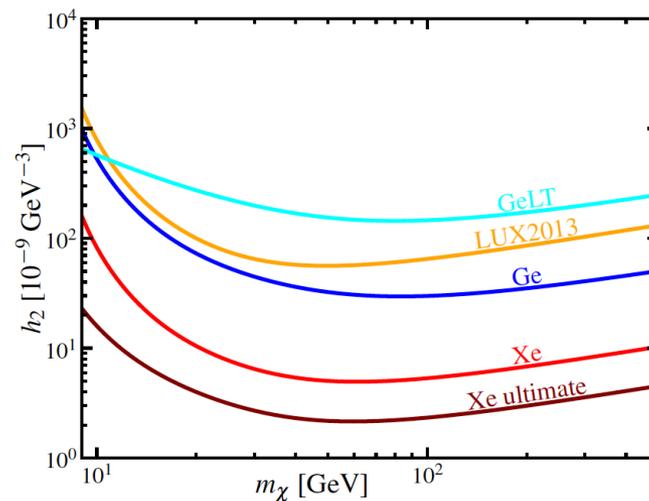
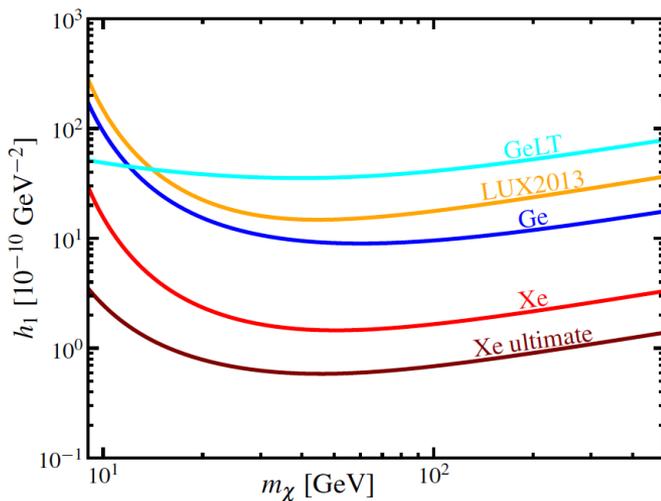
Example case: 20 GeV WIMP with the standard interaction on Xe target, and maximal interaction cross-section.

Conclusions

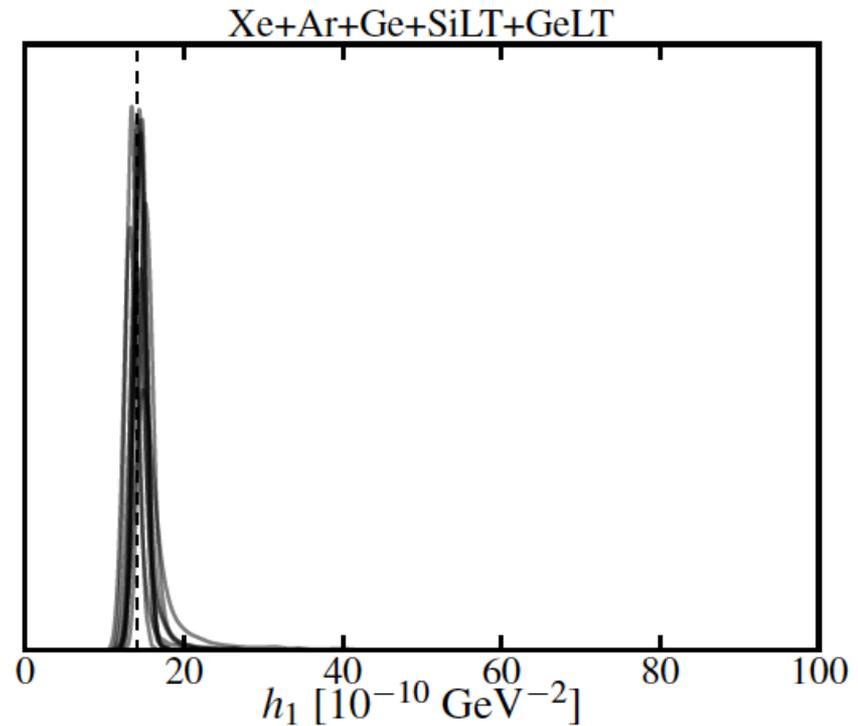
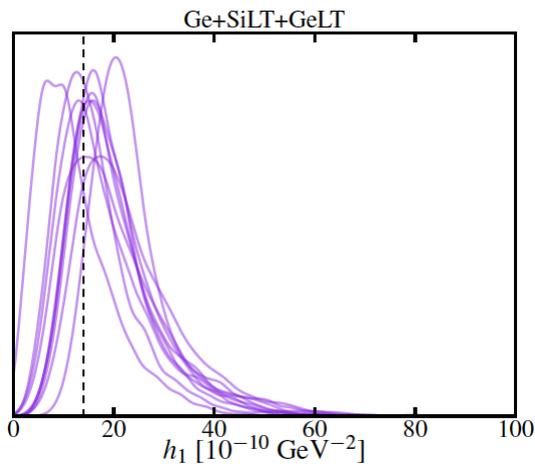
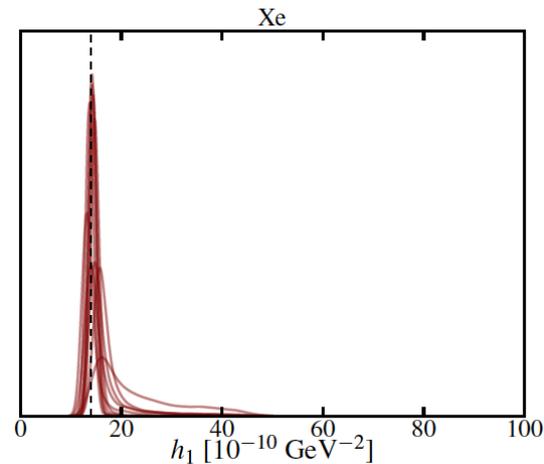
- ✓ A variety of effective operators can describe DM-baryon scattering. We study how distinguishable they are with incoming data.
- ✓ Prospects for model-selection in the context of DM EFT depend on WIMP mass, type of interaction, and experimental capabilities. Low energy thresholds, wide energy window, and large exposures are good.
- ✓ Mass estimates can be biased if data is analyzed assuming a wrong scattering operator. Solution: model selection prior to parameter estimation, or agnostic analysis.
- ✓ Prospects for correctly distinguishing the right operator before reaching the neutrino floor are excellent if signal is just below the current limit, but slim if G2 does not see a signal.



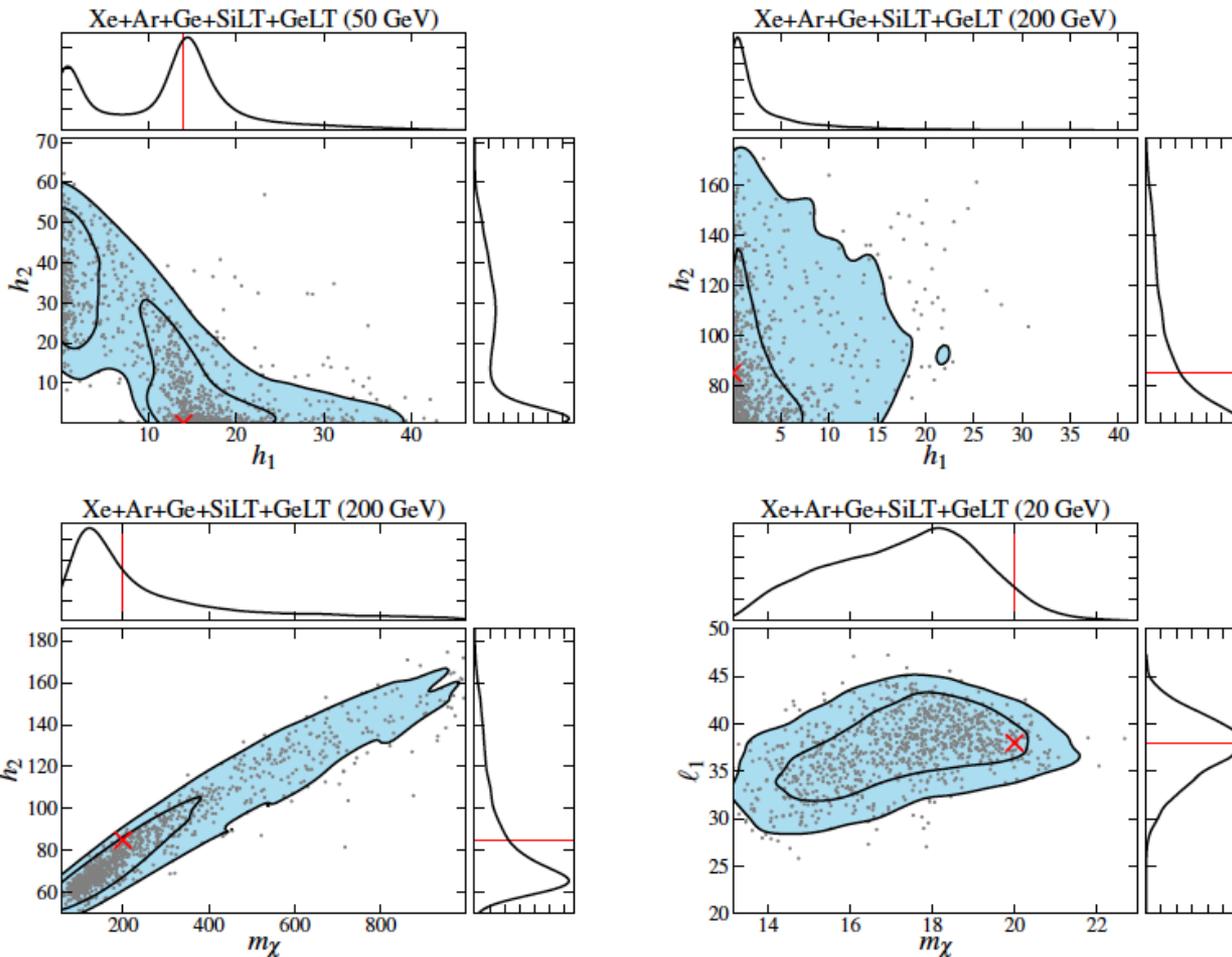
Projected exclusions

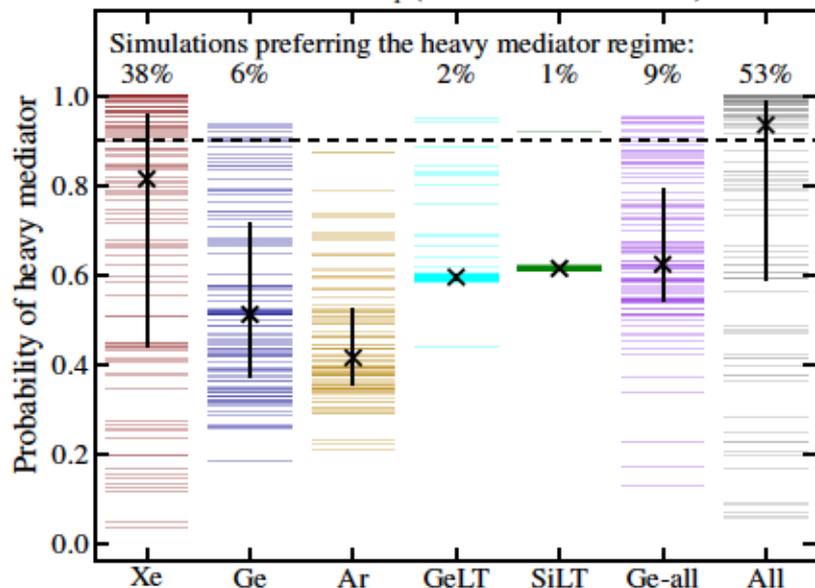
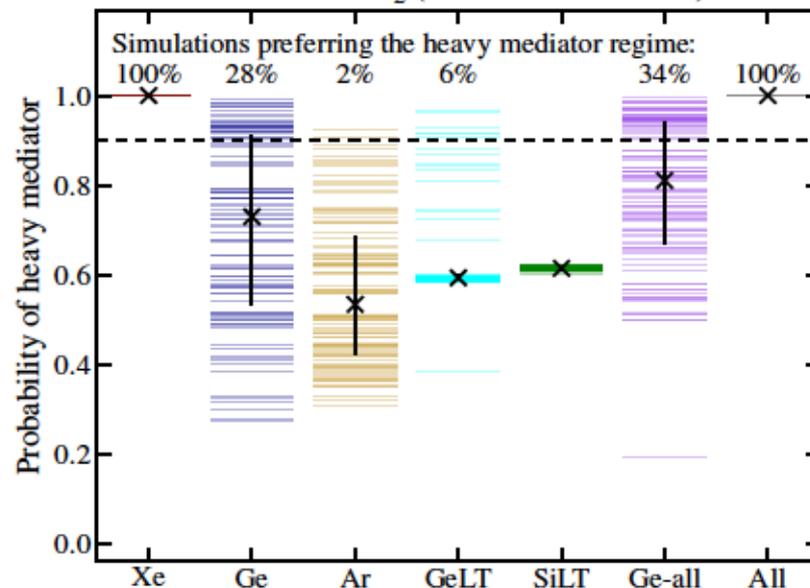
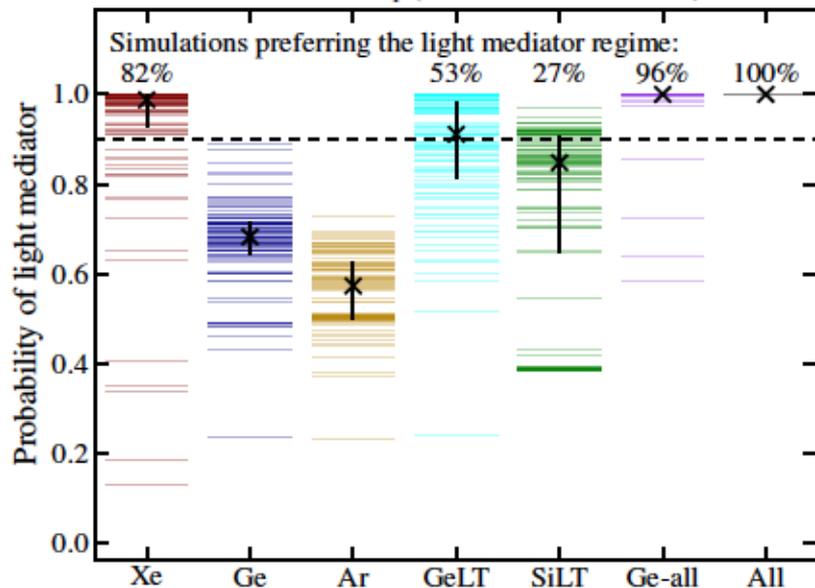
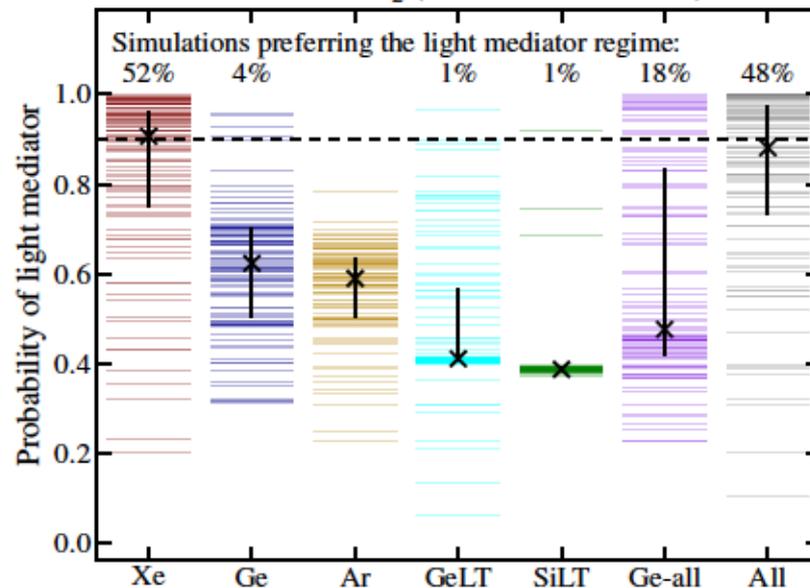


Parameter estimation when true model is known: need large exposure



Multimodal posteriors and large degeneracies



True model: h_1 (WIMP mass=50GeV)True model: h_2 (WIMP mass=50GeV)True model: ℓ_1 (WIMP mass=50GeV)True model: ℓ_2 (WIMP mass=50GeV)

m_χ	$h_1 [10^{-10} \text{ GeV}^{-2}]$	$h_2 [10^{-9} \text{ GeV}^{-3}]$	$\ell_1 [10^{-13}]$	$\ell_2 [10^{-11} \text{ GeV}^{-1}]$
20 GeV	22	100	38	19
50 GeV	14	56	32	14
200 GeV	23	85	54	23

Parameter name	Prior range
m_χ [GeV]	1-1000
h_1 [10^{-10} GeV $^{-2}$]	0.1-10 000
h_2 [10^{-9} GeV $^{-3}$]	0.1-10 000
ℓ_1 [10^{-13}]	0.1-10 000
ℓ_2 [10^{-11} GeV $^{-1}$]	0.1-10 000

WIMP mass	20 GeV	50 GeV	200 GeV
Xe	(88, 91, 87, 90)	(93, 123, 91, 91)	(113, 155, 91, 94)
Ar	(0, 1, 0, 0)	(3, 4, 2, 2)	(7, 10, 3, 5)
Ge	(2, 2, 3, 2)	(2, 3, 3, 3)	(4, 6, 3, 3)
SiLT	(0, 0, 5, 0)	(0, 0, 2, 0)	(0, 0, 1, 0)
GeLT	(0, 0, 27, 2)	(0, 0, 8, 1)	(0, 0, 6, 0)